

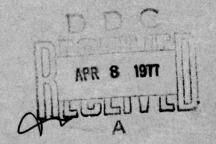
NAVAL OCEANOGRAPHIC OFFICE REFERENCE PUBLICATION 8 - S

WIND DRIFT OF SEA ICE

A SUPPLEMENT TO THE NAVAL OCEANOGRAPHIC OFFICE NUMERICAL ICE FORECASTING SYSTEM

SEPTEMBER 1976





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DEPARTMENT OF THE NAVY WASHINGTON, D. C. 20373

ABSTRACT

The numerical ice prediction system described in RP 8 has been in use since 1971. Its outputs, however, have been related only to ice growth, not ice drift. An addition to the system now provides an output of daily ice drift at 62 arctic stations. Drift is calculated as a function of wind stress, water stress, and Coriolis force. The computer output is sufficiently accurate for ice drift in the locality of the station whose wind observations were input to the model. However, more recent research indicates that for greater accuracy and for computations of divergence and convergence interaction between floes must be considered.

WIND DRIFT OF SEA ICE

A SUPPLEMENT TO THE NAVAL OCEANOGRAPHIC OFFICE NUMERICAL ICE FORECASTING SYSTEM

DONALD J. GERSON LLOYD S. SIMPSON

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FOREWORD

Fleet operations in the Arctic are often hindered, and sometimes aborted, by the severity of ice conditions. The Naval Oceanographic Office has pioneered in the development of ice forecasting support designed to permit the Navy to operate effectively at all times in arctic waters. This publication describes an automated program for computing ice drift for selected locations in the Arctic. It supplements a previously existing program for computing and predicting ice thickness for these locations. Ice drift, as well as ice thickness, information is required for mine warfare, submarine surfacings, surface ship transits, and sonobuoy plants.

J. E. AYRES Captain, USN Commander

CONTENTS

I.	INTRODUCTION	1
II.	DATA	1
III.	METHOD	1
	A. Mean Wind Speed and Direction	1
	B. Drift Computation	3
IV.	VARIATION FOR CONCENTRATION AND ROUGHNESS	9
٧.	CONCLUSIONS	9
VI.	DISCUSSION	9
REFER	RENCES	12
	FIGURES	
1.	Sample output of wind velocity and ice drift	2
2.	Flow chart of ice drift subroutine	4
3a.	Velocity diagram	16
3b.	Force diagram	16
	TABLES	
		10
1.	Factor Table	
	APPENDIXES	
Α.	Subroutine Listing	1:
B.	Theoretical Development of Wind Drift Equations	10

I. INTRODUCTION

Since July 1971 a numerical ice prediction system has been operated by the Naval Oceanographic Office to compute various sea ice characteristics and forecasting aids of importance to arctic operations. Gerson (1975) completely describes the system as of 1975; therefore, it will only be briefly discussed here.

The original outputs of the prediction system include the following:

- * present degree-day accumulations
- * forecast degree-day accumulations
- * daily synoptic temperatures
- * temperature trends
- * selected sea surface temperatures
- * initial freezeup forecasts
- * reported snow depths
- * ice thickness forecasts

All this information relates to ice formation, growth, thickness, and decay in-situ without consideration of ice drift. Since the knowledge of ice drift is extremely important to the forecaster, its inclusion as part of the daily output was considered desirable. The purpose of this report is to describe this significant update to the ice prediction system, namely the computation of wind drift of ice at each of the 62 stations in the present network. An example of a typical output listing is shown in figure 1.

II. DATA

The data are the same as used for the basic system, and are fully described in Gerson (1975). A slight revision in the heading is discussed in U. S. Department of Commerce (1973). The data base remains the entire world's synoptic weather observations collected by the National Oceanic and Atmospheric Administration, but observations for each day now are on one tape rather than two and are processed on its IBM 360/195 computer.

III. METHOD

A. Mean Wind Speed and Direction

The wind drift is derived entirely from values of synoptic wind speed and direction as reported by the network stations. The four (or fewer if reports are missing) daily synoptic wind observations are vectorially averaged to produce a mean wind speed. The equations used are as follows:

$$I = \frac{1}{k} \sum_{i=1}^{k} V_{i}^{sinD}_{i} ,$$

$$J = \frac{1}{k} \sum_{i=1}^{k} v_i^{cosD_i} ,$$

$$\bar{D} = \arctan (I/J)$$

$$\overline{V} = (I^2 + J^2)^{\frac{1}{2}}$$

where

k = number of observations (1-4) during the day,

D; = wind direction at time i,

 V_{i} = wind speed at time i,

 \bar{D} = mean wind direction for the day, and

y = mean wind speed for the day.

The computer program accounts for the quadrant in these computations. This subroutine is listed in appendix A, and a flow chart is shown in figure 2.

B. Drift Computation

The method used for ice drift computation is that of Shuleikin (1953). This method has been used by Wittmann and MacDowell (1964) to produce a set of curves for manual computation of ice drift. In order to obtain a daily output for 62 stations, however, an automated method is necessary. The equations as presented by Shuleikin do not immediately lend themselves to computer programming, and considerable transformation was required. A complete derivation of the equations used in the computer program is given in appendix B. The equations used are as follows:

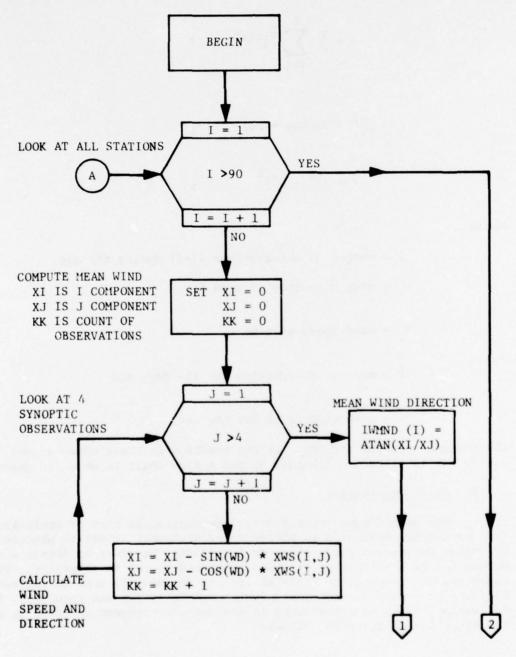
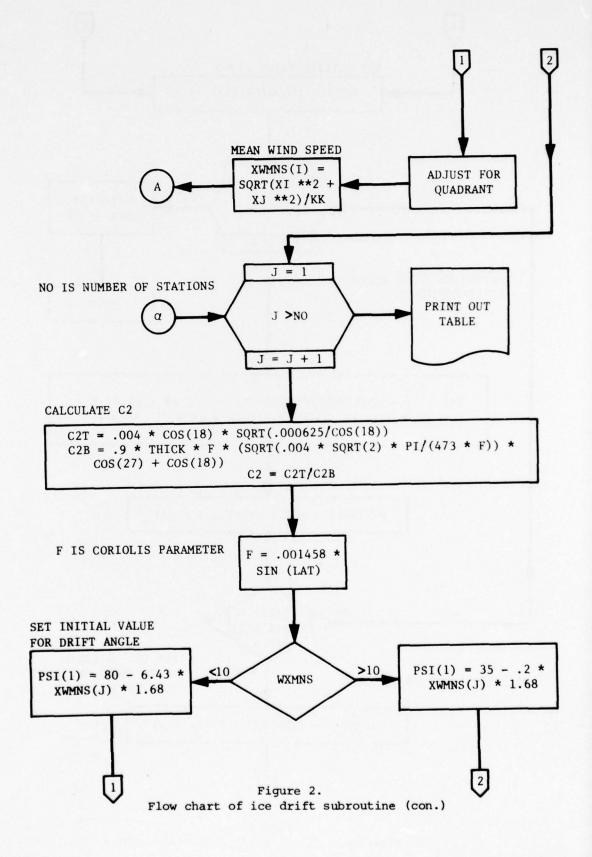


Figure 2. Flow chart of ice drift subroutine



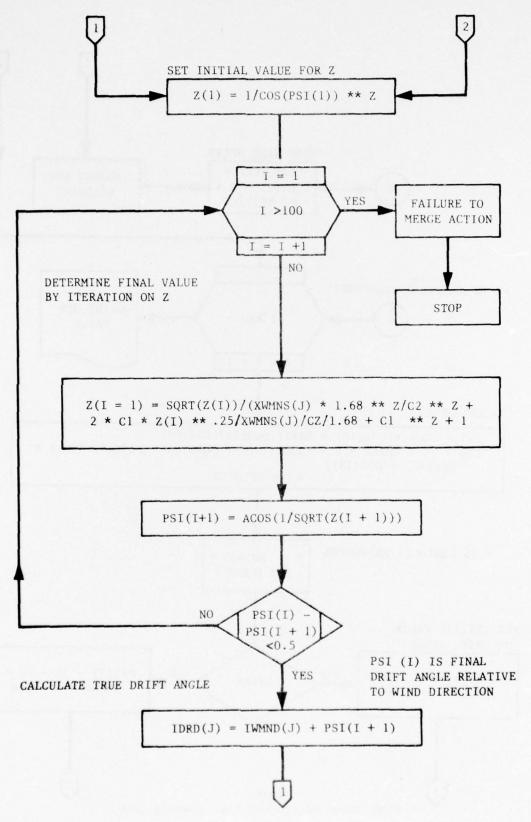


Figure 2. Flow chart of ice drift subroutine (con.)

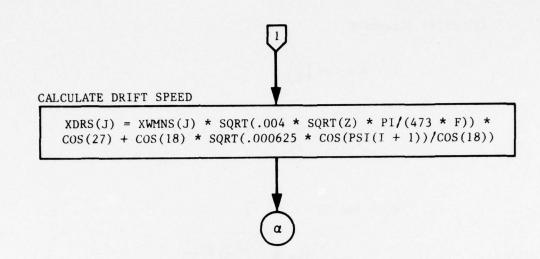


Figure 2. Flow chart of ice drift subroutine (con.)

(1) drift direction

$$\psi = \operatorname{arc} \cos \left(\frac{1}{z}\right)$$
 ,

$$Z = \frac{Z^{\frac{1}{2}}}{V^{2}C_{2}^{2}} + \frac{2C_{1}}{VC_{2}} \quad Z^{\frac{1}{2}} + C_{1}^{2} + 1 \quad ,$$

$$C_1 = \tan 18^{\circ}$$

$$C_{2} = \frac{.004\cos 18^{\circ} \sqrt{\frac{6.25 \times 10^{-4}}{\cos 18^{\circ}}}}{0.9fI \left[\sqrt{\frac{.004 \sqrt{2\pi}}{473f}}\cos 27^{\circ} + \cos 18^{\circ}\right]}$$

$$f = 1.458 \times 10^{-4} \sin(1at)$$
; and

(2) drift speed

$$u = V \left[\sqrt{\frac{.004 \sqrt{2\pi}}{473 f}} \cos 27^{\circ} + \cos 18^{\circ} \right] \sqrt{\frac{6.25 \times 10^{-4} \cos \psi}{\cos 18^{\circ}}}$$
;

where

V = wind speed (ft/sec),

 ψ = drift to the right of the wind direction

u = drift speed (ft/sec)

lat = latitude, and

I = ice thickness (inches)

The equation for Z is required to be in an iterative form, thus it is used as follows:

$$z_{n+1} = \frac{z_n^{\frac{1}{2}}}{\sqrt{c_2^2}} + \frac{2c_1}{vc_2} z_n^{\frac{1}{4}} + c_1^2 + 1$$

IV. VARIATION FOR CONCENTRATION AND ROUGNHESS

The Shuleikin equations assume a single ice floe and do not consider surface roughness. An empirical table of wind factors (Zubov, 1945) was compiled by P.A. Gordienko from many observations of ice drift in the Soviet Arctic. Comparison of this table with the output of the Shuleikin equations shows a match at 0.5 oktas and 4 tenths extent of ridging and hummocking. A factor table (table 1) was then constructed for use in correcting the computer output to the applicable concentration and roughness. Instructions for using the table are given beneath the table. As an example, suppose we have a mean wind speed and direction of 9.2 kn and 102 deg, which produce an ice drift and direction of .34 km and 123 deg (see station 70086 in figure 1). If the concentration were 0.5 oktas and the extent of ridging and hummocking 4 tenths, then, referring to table 1, the ice drift speed would remain .34 kn since the multiplicative factor is 1. If, however, the concentration was 6 oktas and the ridging 3 tenths, the ice drift speed would be .10 km since the multiplicative factor is .3. In either case, the drift direction remains 123 deg, since this table only relates to drift speed.

V. CONCLUSIONS

The methods described above have been programmed and are now operational on the UNIVAC 1108 computer at the Naval Oceanographic Office. Ice drift information is now regularly produced, together with the other products described in the Introduction. The outputs have been evaluated by comparing them first with the manual computation method of Wittmann and MacDowell (1964) and second with manual calculations accomplished with the original equations. These comparisons have proved the computer program to be reliable. Evaluation with actual observed drift conditions has not been accomplished, but since the curves from Wittmann and MacDowell have been used successfully for years, the computer output is assumed accurate.

VI. DISCUSSION

As previously stated, this computation of ice drift is based on Shuleikin's work of 1953. This model was chosen since it had been used manually for years and had been found to give accurate results. In addition, it was most easily adapted to estimating wind drift from station observations that were available in the existing system.

A basic problem exists with the Shuleikin model in that it only considers the motion of a single floe as it results from wind and water stress and Coriolis force. No consideration is given to the effect of interaction between floes. In the program herein described, the effects of concentration and ridging are considered empirically but are not input to the theoretical model. Considerable research has been done since Shuleikin's work. Many of the recent papers are published in the extensive AIDJEX series (University of Washington, 1972-). Campbell and Rasmussen (1972) begin with an excellent

Table 1
FACTOR TABLE

TOTAL CONCENTRATION OF ICE (OKTAS)

	0.5	1	2	3	4	5	6	7.0	7.5
E*									
1	.26	.24	.20	.17	.13	.10	.09	.07	.07
2	.51	.47	.40	.34	.28	.22	.18	.15	.12
3	.75	.71	.63	.55	.47	.39	.30	.22	.19
4	1.00	.96	.84	.72	.62	.51	. 39	.29	.24
5	1.26	1.19	1.06	.92	. 79	.65	.52	.38	. 31
6	1.51	1.43	1.26	1.10	.94	.80	.63	.46	.38
7	1.75	1.66	1.47	1.29	1.10	.92	.73	.55	.46
8	2.01	1.90	1.69	1.49	1.28	1.07	. 85	.64	.53
9	2.26	2.13	1.90	1.64	1.40	1.17	.94	.71	.60

*Extent of ridging and hummocking (tenths)

Table of multiplicative factors for estimating speed of ice drift from computer printout. The drift speeds given in the printout are true for 0.5 oktas and 4 tenths ridging. To determine correct speed enter concentration and extent of ridging into this table and multiply computer's drift speed by the factor thus obtained.

summary of ice dynamics up to 1972 and proceed to give a model that considers ice as a viscous material. Rothrock (1975) considers sea ice as a plastic material. Additional papers are expected shortly.

Ice drift estimates, computed by the model discussed in sections III and IV of this paper, are sufficiently accurate for use in the locality of the station whose wind reports are used in the calculation. They are not sufficiently accurate for the calculation of areas of convergence and divergence or for interpolation between stations or extrapolation into areas lacking wind observations. To add these features to the system in the future, it will be necessary to implement one of the more sophisticated models, requiring more specific inputs not now readily available.

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APPENDIX B. THEORETICAL DEVELOPMENT OF WIND DRIFT EQUATIONS

The frictional stress between two media is assumed to be proportional to the square of the relative velocity between them. Let \vec{V}_a be the wind velocity, \vec{V}_i be the ice velocity, assumed to move in +x direction and \vec{V}_w be the water velocity. The frictional stress (\vec{T}_a) between air and ice is then given by:

$$|\overset{\rightarrow}{\mathbf{T}_{\mathbf{a}}}| = \mathbf{k}_{\mathbf{a}} \rho_{\mathbf{a}} |\overset{\rightarrow}{\mathbf{v}_{\mathbf{a}} - \mathbf{v}_{\mathbf{i}}}|^2 = \mathbf{k}_{\mathbf{a}} \rho_{\mathbf{a}} |\overset{\rightarrow}{\mathbf{v}_{\mathbf{a}}}|^2 \text{ since } |\overset{\rightarrow}{\mathbf{v}_{\mathbf{a}}}| >> |\overset{\rightarrow}{\mathbf{v}_{\mathbf{i}}}|$$

The stress (T_i) between ice and water is given by:

$$|\overrightarrow{\mathbf{T}}_{\mathbf{i}}| = \mathbf{k}_{\mathbf{w}} \rho_{\mathbf{w}} | \overrightarrow{\mathbf{v}}_{\mathbf{i}} - \overrightarrow{\mathbf{v}}_{\mathbf{w}} |^{2}$$

In addition to these two forces we have the Coriolis force:

$$\vec{C} = f(\vec{V}, xk)$$

acting on the ice where $f=2\omega\sin\theta$. The motion is assumed to take place under the equilibrium of these three forces therefore they must vectorially add to zero.

We now choose a coordinate system whose X - axis coincides with the direction of V thus producing figure 3.

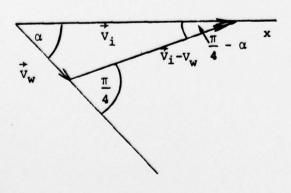


Figure 3a Velocity diagram

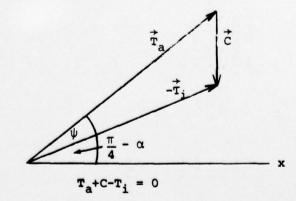


Figure 3b Force diagram

Referring to figure 3a we see:

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 $\frac{\mathbf{w}}{|\vec{\mathbf{v}}_{\mathbf{i}} - \vec{\mathbf{v}}_{\mathbf{w}}|} = \frac{\sin(\pi/4 - \alpha)}{\sin(\alpha)} = \frac{\sin(\pi/4 - \alpha)}{\sin(\pi/4 - \alpha)} = \frac{\sin$ $Ex \text{ forces} = -k_{\text{o}} \rho_{\text{o}} | \vec{v}_{\text{i}} - \vec{v}_{\text{o}} |^{2} \cos(\pi/4 - \alpha) | \vec{T}_{\text{o}} | \cos \theta = 0$ = $\sin(\pi/4)\cos(\alpha)-\sin(\alpha)\cos(\pi/4)$

Solving for sin' and cosy and taking the ratio: (a) nie

= $\sin(\pi/4)\cot(\alpha)-\cos(\pi/4)$ = ζ (0-4\m) $\sin^2\left[\sqrt{y}-\sqrt{y}\right]_{x}$ Q_{x} $x+\left[\sqrt{y}\right]_{x}$ $\cot(\alpha) - \cot(\pi/4) = \cos(\alpha) - 1 = \frac{\zeta}{2}$ k.o. | V. -V. | 2 cos (1/4-0)

 $\cot(\alpha) = 1 + 2\zeta$

again from figure 3a we have:

 $\begin{vmatrix} \overrightarrow{v}_{i} \end{vmatrix} = \begin{vmatrix} \overrightarrow{v}_{i} \end{vmatrix} \cos(\alpha) + \begin{vmatrix} \overrightarrow{v}_{i} \end{vmatrix} - \begin{vmatrix} \overrightarrow{v}_{i} \end{vmatrix} \cos(\pi/4 - \alpha)$

Now referring to figure 3b:

$$T_{ix} = T_i \stackrel{\rightarrow}{\cdot} i = k_w \rho_w | V_i - V_w |^2 \cos(\pi/4 - \alpha)$$

$$T_{iv} = \overrightarrow{T}_{i} \cdot \overrightarrow{j} = k_{iv} \rho_{iv} | \overrightarrow{V}_{i} - \overrightarrow{V}_{iv} |^{2} \sin(\pi/4 - \alpha)$$

$$T_{ax} = \overrightarrow{T}_{v} \cdot \overrightarrow{i} = k_a \rho_a |\overrightarrow{V}_a|^2 \cos \Psi$$

$$T_{av} = \overrightarrow{T} \cdot \overrightarrow{j} = k_a \rho_a |V_a|^2 \sin \Psi$$

$$C_{\mathbf{x}} = \overset{\rightarrow}{\mathbf{C}} \overset{\rightarrow}{\mathbf{i}} = -\mathrm{mf}(\mathbf{k} \mathbf{x} \mathbf{V}_{\mathbf{i}}) \overset{\rightarrow}{\mathbf{i}} = 0 = -\mathrm{mf}(\mathbf{V}_{\mathbf{i}} \mathbf{x} \mathbf{k}) \overset{\rightarrow}{\mathbf{i}}$$

$$C_{\mathbf{y}} = \overset{\rightarrow}{\mathbf{C}} \overset{\rightarrow}{\cdot} \overset{\rightarrow}{\mathbf{J}} = -\mathbf{m} f(\overset{\rightarrow}{\mathbf{k}} \mathbf{x} \overset{\rightarrow}{\mathbf{v}} \overset{\rightarrow}{\mathbf{j}}) \overset{\rightarrow}{\cdot} \overset{\rightarrow}{\mathbf{j}} = -\mathbf{m} f(\overset{\rightarrow}{\mathbf{v}} \overset{\rightarrow}{\mathbf{k}} \overset{\rightarrow}{\cdot} \overset{\rightarrow}{\mathbf{j}})$$

Now since: $|\vec{\mathbf{v}}_{\underline{i}}|$ $|\vec{\mathbf{v}}_{\underline{i}} - \vec{\mathbf{v}}_{\underline{w}}| = \frac{|\vec{\mathbf{v}}_{\underline{i}}|}{\zeta \cos(\alpha) + \cos(\pi/4 - \alpha)}$

and therefore:

$$k_{w}\rho_{w}|\dot{v}_{1}-\dot{v}_{w}|^{2} = \frac{k_{w}\rho_{w}|\dot{v}_{1}|^{2}}{\left[\zeta\cos\left(\alpha\right)+\cos\left(\pi/4-\alpha\right)\right]^{2}}$$

we have:

$$\tan \Psi = \tan (\pi/4-\alpha) + \frac{B}{|\dot{\nabla}_{\dot{1}}|}$$

where:

where: m = mass per unit area

$$B = mf[\cos\alpha + \cos(\pi/4 - \alpha)]^2$$

 $m = \max \text{ per unit area} \qquad \frac{2(\alpha + 4 - \alpha)^2}{m \cdot (\alpha + 4 - \alpha)} = g$ k = the coefficient of friction between ice and water

k, = the coefficient of friction between air and ice

$$T_{\mathbf{v}\mathbf{x}} = |\vec{\mathbf{T}}_{\mathbf{v}}| \cos \Psi = T_{\mathbf{i}\mathbf{x}} = K_{\mathbf{w}} \rho_{\mathbf{w}} |\vec{\mathbf{v}}_{\mathbf{i}} - \vec{\mathbf{v}}_{\mathbf{w}}|^{2} \cos (\pi/4 - \alpha)$$

 ρ_{\bullet} = the density of air

i, j, k, = the unit vectors along the x, y, and z directions respectively

therefore:

$$\Sigma Y \text{ forces} = -k_{\mathbf{w}} \rho_{\mathbf{w}} | \overrightarrow{\mathbf{v}}_{\mathbf{i}} - \overrightarrow{\mathbf{v}}_{\mathbf{w}} |^{2} \sin(\pi/4 - \alpha) - \mathbf{m} \mathbf{f} | \overrightarrow{\mathbf{v}}_{\mathbf{i}} | + | \overrightarrow{\mathbf{T}}_{\mathbf{v}} | \sin \Psi = 0$$

$$\Sigma x \text{ forces } = -k_w \rho_w | \overrightarrow{v}_1 - \overrightarrow{v}_w |^2 \cos(\pi/4 - \alpha) | \overrightarrow{T}_v | \cos \Psi = 0$$

Solving for siny and cosy and taking the ratio:

$$\tan \Psi = \frac{\operatorname{mf} |\overrightarrow{v}_{i}| + k_{w} \rho_{w} |\overrightarrow{v}_{i} - \overrightarrow{v}_{w}|^{2} \sin(\pi/4 - \alpha)}{k_{w} \rho_{w} |\overrightarrow{v}_{i} - \overrightarrow{v}_{w}|^{2} \cos(\pi/4 - \alpha)}$$

$$= \tan(\pi/4 - \alpha) + \frac{\operatorname{mf} |\overrightarrow{v}_{i}|}{k_{w} \rho_{w} |\overrightarrow{v}_{i} \overrightarrow{v}_{w}|^{2} \cos(\pi/4 - \alpha)}$$

Now since:

$$|\overrightarrow{v}_{i} - \overrightarrow{v}_{w}| = \frac{|\overrightarrow{v}_{i}|}{\zeta \cos(\alpha) + \cos(\pi/4 - \alpha)}$$

and therefore:

$$k_{\mathbf{w}} \rho_{\mathbf{w}} | \mathbf{v}_{\mathbf{i}}^{\dagger} - \mathbf{v}_{\mathbf{w}}^{\dagger} |^{2} = \frac{k_{\mathbf{w}} \rho_{\mathbf{w}} | \mathbf{v}_{\mathbf{i}}^{\dagger} |^{2}}{\left[\zeta \cos{(\alpha)} + \cos{(\pi/4 - \alpha)} \right]^{2}}$$

we have:

$$\tan \Psi = \tan (\pi/4 - \alpha) + \frac{B}{|\vec{v}_i|}$$

where:

$$B = \frac{mf[\zeta\cos\alpha+\cos(\pi/4-\alpha)]^2}{k_{\omega}\rho_{\omega}\cos(\pi/4-\alpha)}$$

Now since:

$$T_{vx} = |\vec{T}_v| \cos \Psi = T_{ix} = k_w \rho_w |\vec{v}_i - \vec{v}_w|^2 \cos (\pi/4 - \alpha)$$

$$= \frac{k_{\mathbf{w}} \rho_{\mathbf{w}} |V_{\mathbf{i}}|^2 \cos(\pi/4-\alpha)}{\left[\zeta \cos(\alpha) + \cos(\pi/4-\alpha)\right]^2}$$

we get:

$$|\vec{T}_{\mathbf{v}}| = \frac{k_{\mathbf{w}} \rho_{\mathbf{w}} |\vec{\mathbf{v}}_{\mathbf{i}}|^{2} \cos(\pi/4 - \alpha)}{\left[\zeta \cos(\alpha) + \cos(\pi/4 - \alpha)\right]^{2} \cos^{\Psi}}$$

Now since:

$$|\overrightarrow{\mathbf{T}}_{\mathbf{v}}| = k_{\mathbf{a}} \rho_{\mathbf{a}} |\overrightarrow{\mathbf{v}}_{\mathbf{a}}|^2$$

we get:

$$k_{\mathbf{a}} \rho_{\mathbf{a}} | \overrightarrow{\mathbf{v}}_{\mathbf{a}} | = \frac{k_{\mathbf{w}} \rho_{\mathbf{w}} | \overrightarrow{\mathbf{v}}_{\mathbf{i}} |^{2} \cos(\pi/4 - \alpha)}{\left[\zeta \cos(\alpha) + \cos(\pi/4 - \alpha) \right]^{2} \cos \Psi}$$

or:

$$\frac{|\vec{v}_{i}|}{|\vec{v}_{a}|} = \left[\zeta\cos(\alpha) + \cos(\pi/4 - \alpha)\right] \sqrt{\frac{k_{a}\rho_{a}\cos(\Psi)}{k_{w}\rho_{w}\cos(\pi/4 - \alpha)}}$$

Now from above we have:

$$\frac{B}{|\vec{V}_i|} = \tan(\Psi) - \tan(\pi/4 - \alpha)$$

and multiplying this by the previous equation we get:

$$\frac{B}{|\vec{V}_{a}|} = \left[\zeta\cos(\alpha) + \cos(\pi/4 - \alpha)\right] \sqrt{\frac{k_{a}\rho_{a}\cos\Psi}{k_{w}\rho_{w}\cos(\pi/4 - \alpha)}} \left[\tan(\Psi) - \tan(\pi/4 - \alpha)\right]$$

The equation for ζ is derived from Ekman theory, as follows:

$$v_{w} = \frac{k_{w}\gamma^{2}}{\sqrt{f}} \sqrt{\frac{\rho_{w}}{\mu}}$$

where: $\gamma = |\overrightarrow{v}_i - \overrightarrow{v}_w|$

 μ = the coefficient of viscosity between ice and water

$$\sqrt{\frac{\rho_{\mathbf{w}}}{\mu}} = \frac{\pi}{N |\vec{\nabla}_{\mathbf{w}}| \sqrt{f}}$$

Where: $N = D/|\vec{V}_1|$ and is the depth of frictional mixing. This has been evaluated as N = 473, (Shuleikin). Now since:

$$\zeta = \frac{|\vec{\mathbf{v}}_{\underline{\mathbf{i}}}|}{\gamma}$$

we derive:

$$\zeta = \sqrt{\frac{\pi k_{w} \sqrt{2}}{Nf}}$$

$$=\sqrt{\frac{\pi k_{w}\sqrt{2}}{473f}}$$

It can be assumed with sufficient accuracy that the density of air (ρ_a) is 1.25×10^{-3} gm per cm³, the density of water is 1 gm per cm³, the coefficient of friction between ice and water (k_w) is 4×10^{-3} and between air and ice (k_a) is 2×10^{-3} . From observations it has been determined that $\alpha = 27^{\circ}$. In addition, since the density of ice is .9 and the mass per unit area of ice (m) is the density times the thickness (L), m = .9L. Therefore: from the equation for $\frac{B}{|V_a|}$ above, we get:

$$\frac{1}{|\overrightarrow{V}_a|} = \frac{\rho_{\mathbf{w}} k_{\mathbf{w}} \cos(18^\circ)}{.9 \text{Lf}[\zeta \cos(27^\circ) + \cos(18^\circ)]} \sqrt{\frac{6.25 \times 10^{-4} \cos(\Psi)}{\cos(18^\circ)}} \left[\sqrt{\sec^2(\Psi) - 1} - \tan(18^\circ)\right]$$

let:
$$Z = \frac{1}{\cos^2(\Psi)}$$

then
$$\frac{1}{|\vec{v}_a|} = \frac{.004\cos(18^\circ)\frac{1}{2}}{.9Lf[\frac{.004}{473f}} \frac{\sqrt{\frac{6.25\times10^4}{\cos18^\circ}}}{\cos(27^\circ)+\cos(18^\circ)}]$$

$$\tan (18^{\circ}) = \frac{z^{\frac{1}{4}}.9Lf[\frac{.004 2\pi}{473f} \cos 27^{\circ} + \cos 18^{\circ}]}{|\vec{v}_{a}|[.004\cos 18^{\circ} \sqrt{\frac{6.25 \times 10^{-4}}{\cos 18^{\circ}}}]} = \sqrt{z-1}$$

let:
$$C = tan(18^\circ)$$

$$C_{2} = \frac{.004\cos(18^{\circ})\sqrt{\frac{6.25\times10^{-4}}{\cos(18^{\circ})}}}{.9Lf\left[\frac{.004\ 2\pi}{473f}\cos(27^{\circ})+\cos(18^{\circ})\right]}$$

then:
$$\sqrt{Z-1} = C_1 + \frac{z^{\frac{1}{4}}}{C_2 |\vec{\nabla}_a|}$$

$$z = c_1^2 + \frac{z^{\frac{1}{2}}}{c_2^2 |\vec{v}_a|^2} + \frac{2c_1 z^{\frac{1}{4}}}{c_2 |\vec{v}_a|} + 1$$

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The numerical ice prediction system described in RP 8 has been in use since 1971. Its outputs, however, have been related only to ice growth, not ice drift. An addition to the system now provides an output of daily ice drift at 62 arctic stations. Drift is calculated as a function of wind stress, water stress, and Coriolis force. The computer output is sufficiently accurate for ice drift in the locality of the station whose wind observations were input to the model. However, more recent research indicates that for

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SUPPLEMENTARY

INFORMATION

ERRATA

p. 16	figure 3a	change $\vec{v_i}$ - $\vec{v_w}$ to $\vec{v_i}$ - $\vec{v_w}$
p. 17	line 5	" $cos(\alpha)-1$ to $cot(\alpha)-1$
"	" 6	" 1+2; to 1+1/2;
"	" 12	" T _v to T _a
11	" 13	" T _v to T _a
p. 18	" 2	" $ \tilde{T}_{v} $ to $ \tilde{T}_{a} $
"	" 3	$\cos(\pi/4-\alpha) \overset{+}{T}_{v} \text{to } \cos(\pi/4-\alpha) + \overset{+}{T}_{a} $
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p. 19	" 3	" $ \mathring{T}_{v} $ to $ \mathring{T}_{a} $
"	" 5	" Tv to Ta
p. 21	" 2	" $\sqrt{\frac{6.25 \times 10^4}{\cos 18^6}}$ to $\sqrt{\frac{6.25 \times 10^{-4}}{\cos 18^6}}$
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